

Vectorized operator evaluations for solutions of PDEs on heterogeneous architectures with libCEED

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CEED
EXASCALE DISCRETIZATIONS


EXASCALE COMPUTING PROJECT

Where do I come from?



Siena, Tuscany, Italy



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Ok, where do I come from, mathematically speaking?

Master's thesis:
Catmull-Clark *Subdivision Surfaces* (Computer-Aided Design, '78)

The limit-surface obtained is a bicubic uniform B-spline surface of class C^2 , except at *extraordinary* points



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PhD dissertation

PhD in Applied Math from NJIT on numerical simulations of thin films
(long-waves) of viscoelastic fluids

Second order finite
differences on a
staggered grid

Fixed Δx ,
adaptive Δt
Crank-Nicolson



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PhD dissertation (cont'ed)

Long-wave approximation in two spatial dimensions:

$$(1 + \lambda_2 \partial_t) h_t + \frac{\partial}{\partial x} \left\{ (\lambda_2 - \lambda_1) \left(\frac{h^2}{2} Q - h R \right) h_t + \left[(1 + \lambda_1 \partial_t) \frac{h^3}{3} + (1 + \lambda_2 \partial_t) b h^2 \right] \frac{\partial}{\partial x} \left(\frac{\partial^2 h}{\partial x^2} + \Pi(h) \right) \right\} = 0,$$

$$Q + \lambda_2 Q_t = -\frac{\partial}{\partial x} \left(\frac{\partial^2 h}{\partial x^2} + \Pi(h) \right),$$

$$R + \lambda_2 R_t = -h \frac{\partial}{\partial x} \left(\frac{\partial^2 h}{\partial x^2} + \Pi(h) \right).$$

$$\text{disjoining pressure: } \Pi(h) = \frac{\gamma(1-\cos\theta_e)}{M h_*} \left[\left(\frac{h_*}{h} \right)^n - \left(\frac{h_*}{h} \right)^m \right],$$

θ_e contact angle, $M = 0.5$, ($n = 3$, $m = 2$), h_* precursor film thickness.

Jeffreys' constitutive law:

$$\sigma + \lambda_1 \partial_t \sigma = 2\eta(\dot{\epsilon} + \lambda_2 \partial_t \dot{\epsilon})$$



PhD dissertation (cont'ed)

Shear and extensional free-boundary flows of viscoelastic membranes

Linear finite elements with plane stress formulation. Different constitutive models considered: elastic, viscous, viscoelastic (Maxwell)



References

Some references:

V. Barra, S. AFKHAM, L. KONDIC, *Mathematical and numerical modeling of thin viscoelastic films of Jeffreys type subjected to the van der Waals and gravitational forces*, EPJE, **42**, 1 – 14 (2019)

V. Barra, S. A. CHESTER, S. AFKHAM, *Numerical Simulations of Nearly Incompressible Viscoelastic Membranes*, Computers & Fluids, **175**, 36 – 47 (2018)

V. Barra, S. AFKHAM, L. KONDIC, *Interfacial dynamics of thin viscoelastic films and drops*, J. Non-Newt. Fluid Mech. **237**, 26 – 38 (2016)



Internship at Pixar

12-week internship program at



Developed 2 proprietary libraries in C++ for viscous fluid simulations on curved surfaces

A vorticity-formulation Navier-Stokes solver with fluid-structure interactions, using Discrete Exterior Calculus (DEC) and time-splitting



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Internship at Pixar

A thin film (long-wave) solver on curved surfaces

Arbitrary topology
and element shapes

Prototyped a plugin
for Side FX Houdini



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Overview

① Introduction

② libCEED

③ Performance

Theta

Skylake

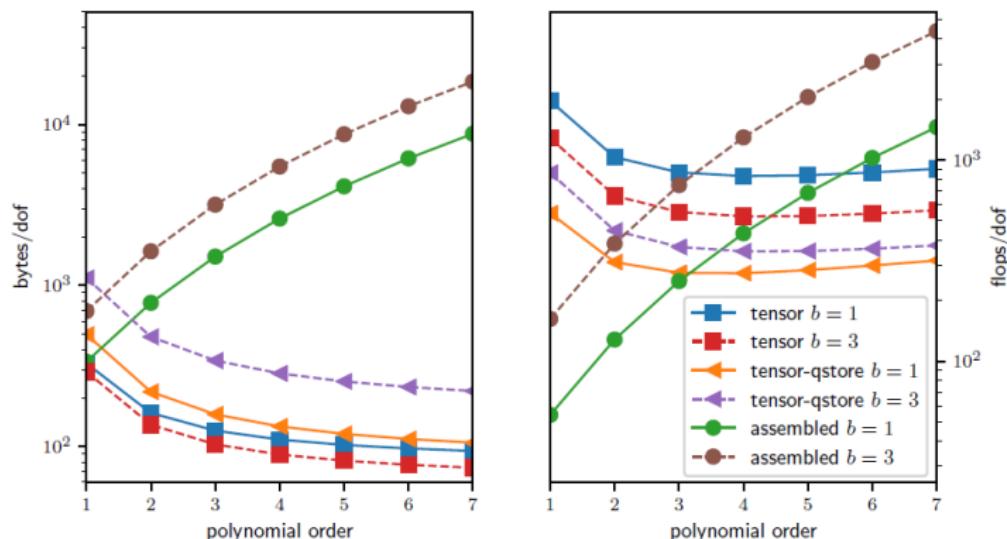
Summit

④ Numerical Examples



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Motivation: Why matrix-free? And why high-order?



Memory bandwidth (left) and FLOPs per dof (right) to apply a Jacobian matrix, obtained from discretizations of a b -variable PDE system. Assembled matrix vs matrix-free (exploits the tensor product structure by either storing at q-points or computing on the fly)

[Courtesy: Jed Brown]

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Overview

- For decades, high-order numerical methods have been considered too expensive
- A sparse matrix is no longer a good representation for high-order operators. In particular, the Jacobian of a nonlinear operator is known to rapidly lose sparsity as the order is increased
- libCEED uses a matrix-free operator description, based on a purely algebraic interface, where user only specifies action of weak form operators
- libCEED operator representation is optimal with respect to the FLOPs needed for its evaluation, as well as the memory transfer needed for operator evaluations (matvec)
 - Matrix-free operators that exploit tensor-product structures reduce the work load from $O(p^6)$ (for sparse matrix) to $O(p^4)$, and memory storage from $O(p^6)$ to $O(p^3)$
- We demonstrate the usage of libCEED, its integration with other packages, and some PETSc application examples



libCEED: the library of CEED

(Center for Efficient Exascale Discretizations)

- Primary target: high-order finite/spectral element methods (FEM/SEM) exploiting tensor-product structure
- Open source (BSD-2 license) C library with Fortran and Python interfaces
- Releases: v0.1 (January 2018), v0.2 (March 2018), v0.3 (September 2018), v0.4 (March 2019), v0.5 (September 2019)



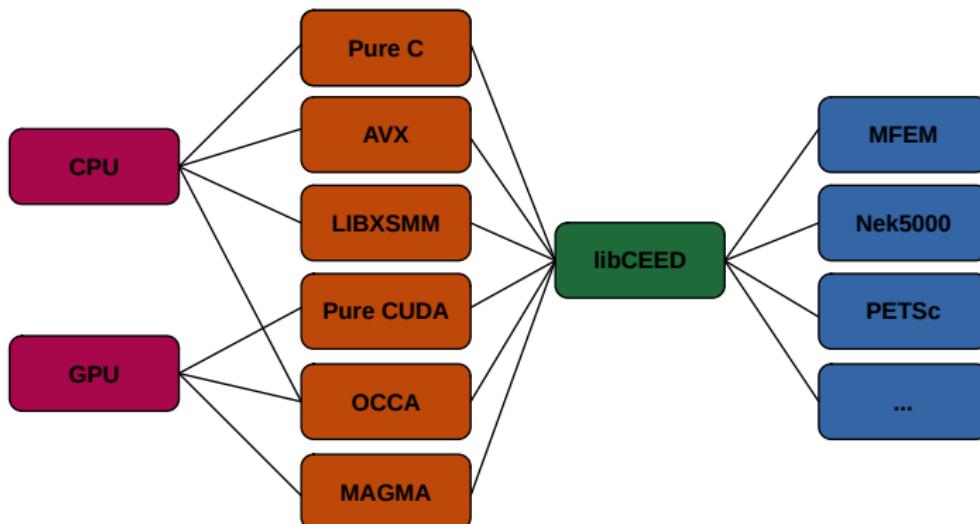
For latest release:

Tomov S., Abdelfattah A., **Barra V.**, Beams N., Brown J. et al., *CEED ECP Milestone Report: Performance tuning of CEED software and 1st and 2nd wave apps* (2019, Oct 2nd) DOI: <https://doi.org/10.5281/zenodo.3477618>



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libCEED backends



libCEED backends

/cpu/self/ref/*:	with * reference serial and blocked implementations
/cpu/self/avx/*:	AVX (Advanced Vector Extensions instruction sets) with * reference serial and blocked implementations
/cpu/self/xsmm/*:	LIBXSMM (Intel library for small dense/sparse mat-multiply) with * reference serial and blocked implementations
/*/occa:	OCCA (just-in-time compilation) with *: CPU, GPU, OpenMP (Open Multi-Processing: API), OpenCL (framework for CPUs, GPUs, etc.)
/gpu/magma:	CUDA MAGMA (dense Linear Algebra library for GPUs and multicore architectures) kernels
/gpu/cuda/*:	CUDA with *: ref (reference pure CUDA kernels), reg (CUDA kernels using one thread per element), shared , optimized CUDA kernels using shared memory gen , optimized CUDA kernels using code generation

Same source code can call multiple CEDDs with different backends. On-device operator implementation with unique interface

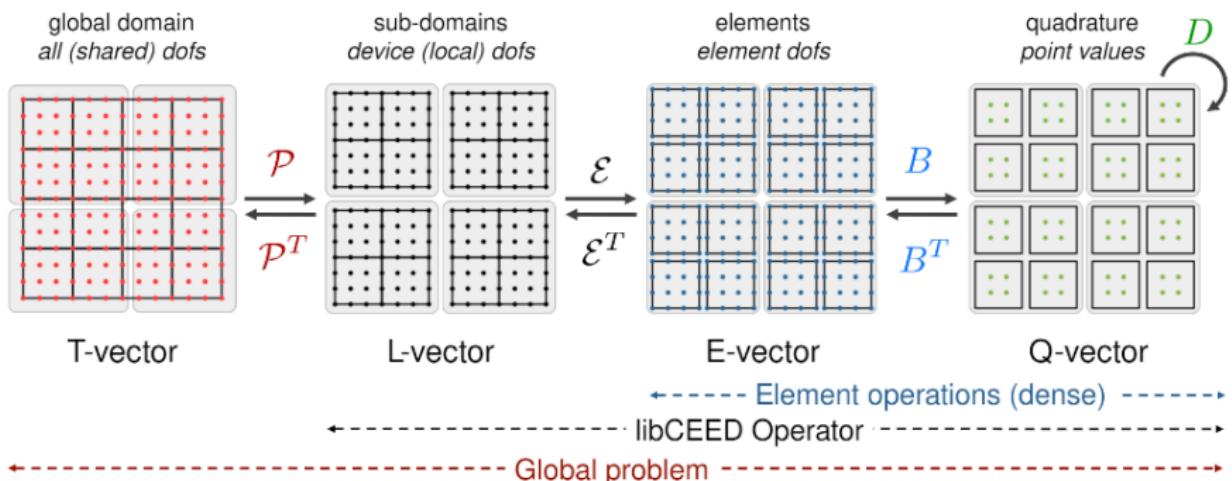


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libCEED decomposition



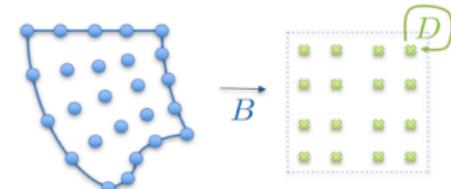
$$A = \mathcal{P}^T \mathcal{E}^T \mathcal{B}^T \mathcal{D} \mathcal{B} \mathcal{E} \mathcal{P}$$



libCEED API objects

- \mathcal{E} : Ceed Element Restriction

Restrict to single element
User choice in ordering



- \mathbf{B} : Ceed Basis Applicator

Describes the actions on basis such as interpolation, gradient, div, curl
Independent of geometry and element topology

- \mathcal{D} : Ceed QFunction

Operator that defines the action of the physics at quadrature points
Choice of interlaced (by fields) or blocked (by element) for multi-component vectors

- $\mathbf{C} = \mathcal{E}^T \mathbf{B}^T \mathcal{D} \mathbf{B} \mathcal{E}$: CeedOperator

Composition of different operators defined on different element topologies possible

- $\mathbf{A} = \mathbf{P}^T \mathbf{C} \mathbf{P}$: User code responsible for parallelization on different compute devices. We use PETSc



Point-wise QFunctions

User-defined
QFunctions:

$$-\nabla \cdot (\kappa(\mathbf{x}) \nabla \mathbf{u}) - \mathbf{f}$$



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Point-wise QFunctions

User-defined
QFunctions:

$$-\nabla \cdot (\kappa(\mathbf{x}) \nabla u) - f$$

or from libCEED's
Gallery:

$$\nabla \cdot (\nabla u) - f$$

are point-wise
functions that do not
depend on element
resolution, topology,
or basis order



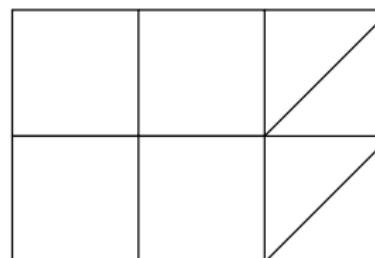
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Point-wise QFunctions

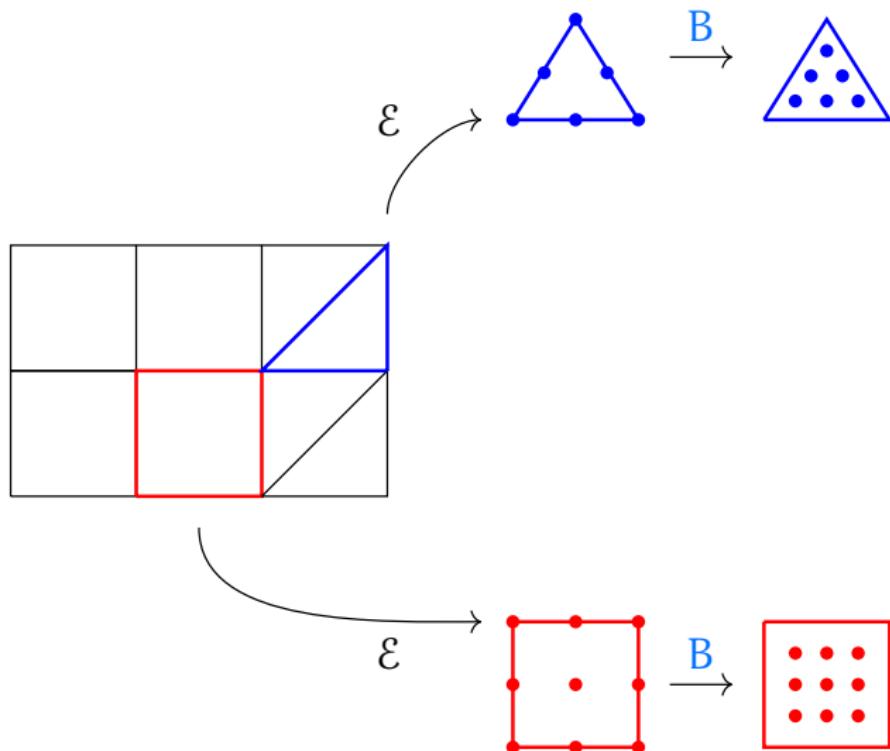
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QFunctions:

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or from libCEED's
Gallery:

$$\nabla \cdot (\nabla u) - f$$

are point-wise
functions that do not
depend on element
resolution, topology,
or basis order



libCEED API for operator composition

Creation of QFunctions and CEED operators:

```
CeedQFunctionCreateInterior(ceed, 1, Mass, Mass_loc, &qf_mass);
CeedQFunctionAddInput(qf_mass, "u", 5, CEED_EVAL_INTERP);
CeedQFunctionAddInput(qf_mass, "weights", 1, CEED_EVAL_NONE);
CeedQFunctionAddOutput(qf_mass, "v", 5, CEED_EVAL_INTERP);

CeedOperatorCreate(ceed, qf_mass, NULL, NULL, &op_mass);
CeedOperatorSetField(op_mass, "u", Erestrictu, CEED_TRANSPOSE, basisu, CEED_VECTOR_ACTIVE);
CeedOperatorSetField(op_mass, "weights", Erestrictudi, CEED_NOTTRANSPOSE, basisx, weights);
CeedOperatorSetField(op_mass, "v", Erestrictu, CEED_TRANSPOSE, basisu, CEED_VECTOR_ACTIVE);
```



libCEED API for operator composition

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```

$$f(\mathbf{u}, \mathbf{p}; \Theta) \rightleftharpoons f(\mathbf{u}, \mathbf{p}; \Theta)$$



libCEED API for operator composition

Creation of QFunctions and CEED operators:

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CeedQFunctionCreateInterior(ceed, 1, Mass, Mass_loc, &qf_mass);
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CeedOperatorSetField(op_mass, "v", Erestrictu, CEED_TRANSPOSE, basisu, CEED_VECTOR_ACTIVE);
```

$$f(\mathbf{u}, \mathbf{p}; \Theta) \rightleftharpoons f(\mathbf{u}, \mathbf{p}; \Theta)$$

Composition of operators for multiphysics or mixed element meshes:

```
CeedCompositeOperatorCreate(ceed, &op_comp);
CeedCompositeOperatorAddSub(op_comp, op_1);
CeedCompositeOperatorAddSub(op_comp, op_2);
```



libCEED's Python interface

Classes:

Ceed

Vector

ElemRestriction

Basis

QFunction

Operator



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libCEED's Python interface

Classes:

Ceed

Vector

ElemRestriction

Basis

QFunction

Operator

CeedVector's data



`numpy.array`



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libCEED's Python interface

Classes:

Ceed

Vector

ElemRestriction

Basis

QFunction

Operator

CeedVector's data



numpy.array

```
import libCEED
ceed = libCEED.Ceed()
n = 10
x = ceed.Vector(n)

print(x)
```

or

```
import libCEED
ceed = libCEED.Ceed()
x = ceed.Vector(size=10)

print(x)
```



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Python interface (some examples)

Vector

```
import libceed
ceed = libceed.Ceed('/gpu/cuda/gen')
x = ceed.Vector(size=10)

a = np.arange(1, 4, dtype="float64")
x.set_array(a, cmode=libceed.USE_POINTER)
b = x.get_array_read()
print(b)
```

```
import libceed
ceed = libceed.Ceed('/cpu/self')
u = ceed.Vector(size=10)

u.set_value(0)
print(u)
```



Python interface (some examples)

Vector

```
import libceed
ceed = libceed.Ceed('/gpu/cuda/gen')
x = ceed.Vector(size=10)

a = np.arange(1, 4, dtype="float64")
x.set_array(a, cmode=libceed.USE_POINTER)
b = x.get_array_read()
print(b)
```

```
import libceed
ceed = libceed.Ceed('/cpu/self')
u = ceed.Vector(size=10)

u.set_value(0)
print(u)
```

Elemrestriction

```
ne = 8
x = ceed.Vector(ne+1)
y = ceed.Vector(2*ne)
r = ceed.ElemRestriction(ne, 2, ne+1, 1, ind, cmode=libceed.USE_POINTER)
r.apply(x, y)
r.T.apply(y, x)
```



C interface Vs Python interface

```
CeedQFunctionCreateInterior(ceed, 1, Mass, Mass_loc, &qf_mass);
CeedQFunctionAddInput(qf_mass, "u", 5, CEED_EVAL_INTERP);
CeedQFunctionAddInput(qf_mass, "weights", 1, CEED_EVAL_NONE);
CeedQFunctionAddOutput(qf_mass, "v", 5, CEED_EVAL_INTERP);

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CeedOperatorSetField(op_mass, "u", Erestrictu, CEED_TRANSPOSE, basisu, CEED_VECTOR_ACTIVE);
CeedOperatorSetField(op_mass, "weights", Erestrictudi, CEED_NOTRANSPOSE,
CEED BASIS_COLLOCATED, weights);
CeedOperatorSetField(op_mass, "v", Erestrictu, CEED_TRANSPOSE, basisu, CEED_VECTOR_ACTIVE);
```

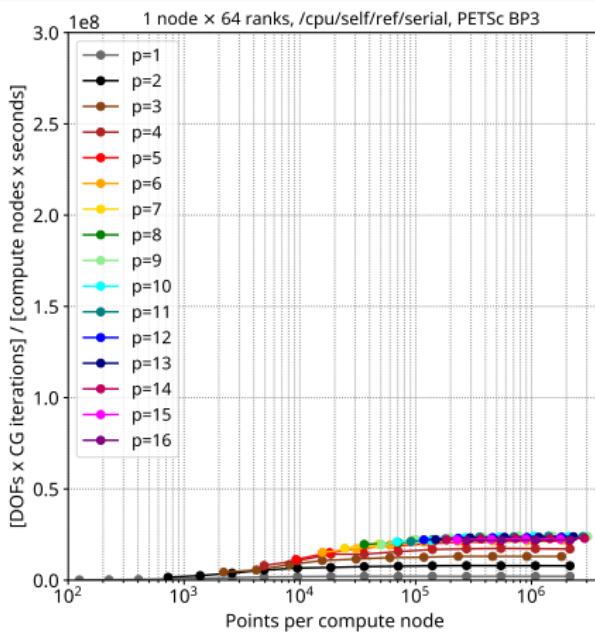
```
qf_mass = ceed.QFunction(1, qfs.apply_mass, os.path.join(file_dir,
"test-qfunctions.h:apply_mass"))
qf_mass.add_input("u", 5, libceed.EVAL_INTERP)
qf_mass.add_input("weights", 1, libceed.EVAL_NONE)
qf_mass.add_output("v", 5, libceed.EVAL_INTERP)

op_mass = ceed.Operator(qf_mass)
op_mass.set_field("u", Erestrictu, basisu, libceed.VECTOR_ACTIVE, libceed.TRANSPOSE)
op_mass.set_field("weights", Erestrictudi, libceed.BASIS_COLLOCATED, weights)
op_mass.set_field("v", Erestrictu, basisu, libceed.VECTOR_ACTIVE, libceed.TRANSPOSE)
```

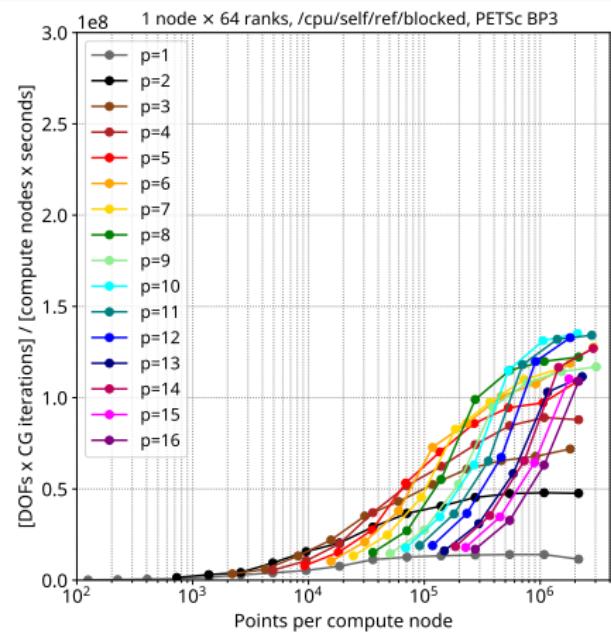


Theta

Performance w. r. t. problem size on KNL: ref



(a)



(b)

Figure: Knight Landing (Intel Xeon Phi 7230 SKU 1.3 GHz) with Intel-18 compiler. In (a) serial implementation; in (b) blocked implementation ($q = P + 2$, $P = p + 1$)



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Performance w. r. t. time on KNL: ref

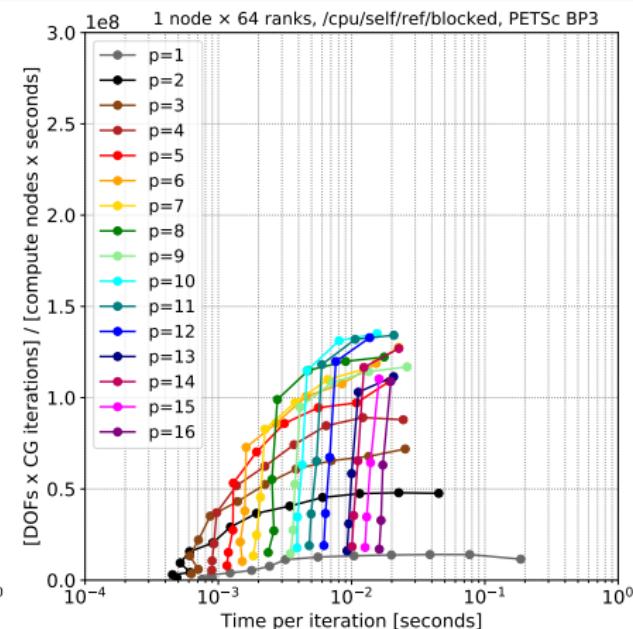
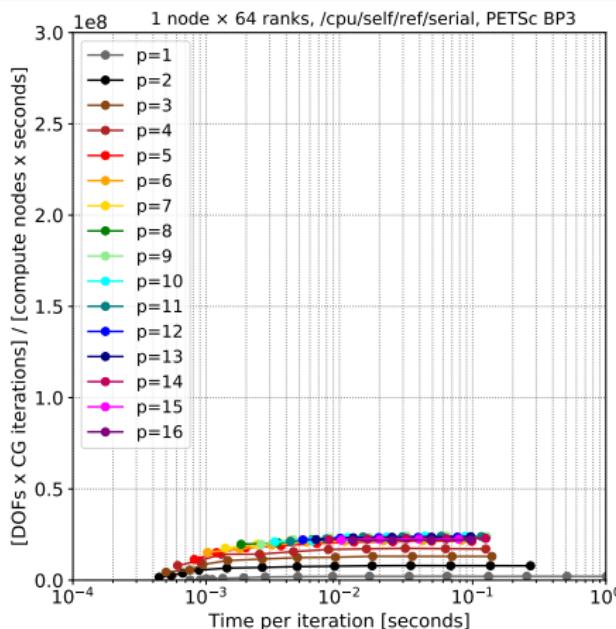
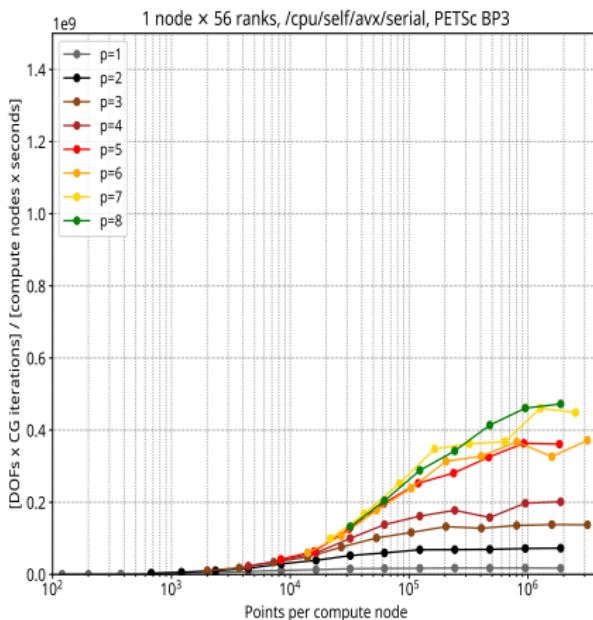
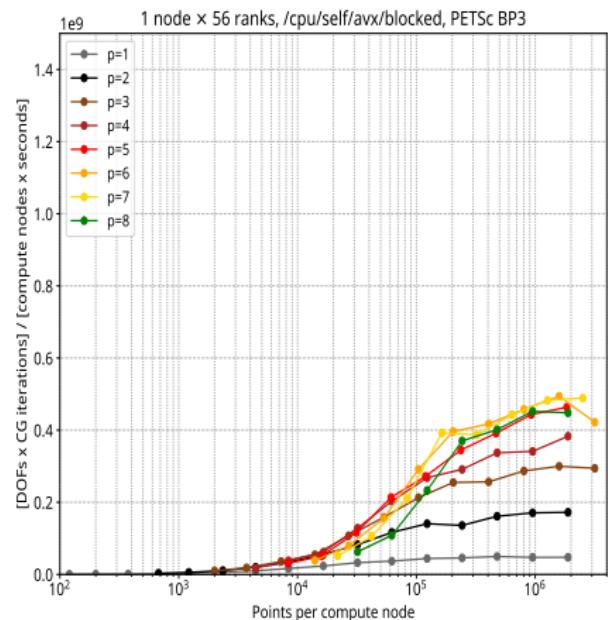


Figure: Knight Landing (Intel Xeon Phi 7230 SKU 1.3 GHz) with Intel-18 compiler. In (a) serial implementation; in (b) blocked implementation ($q = P + 2$, $P = p + 1$)

Performance w. r. t. problem size on Skylake: AVX



(a)



(b)

Figure: Skylake (2x Intel Xeon Platinum 8180M CPU 2.50GHz) with gcc-8 compiler. Backends: in (a) AVX serial; in (b) AVX blocked ($q = P + 2$, $P = p + 1$)



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Skylake

Performance w. r. t. time on Skylake: AVX

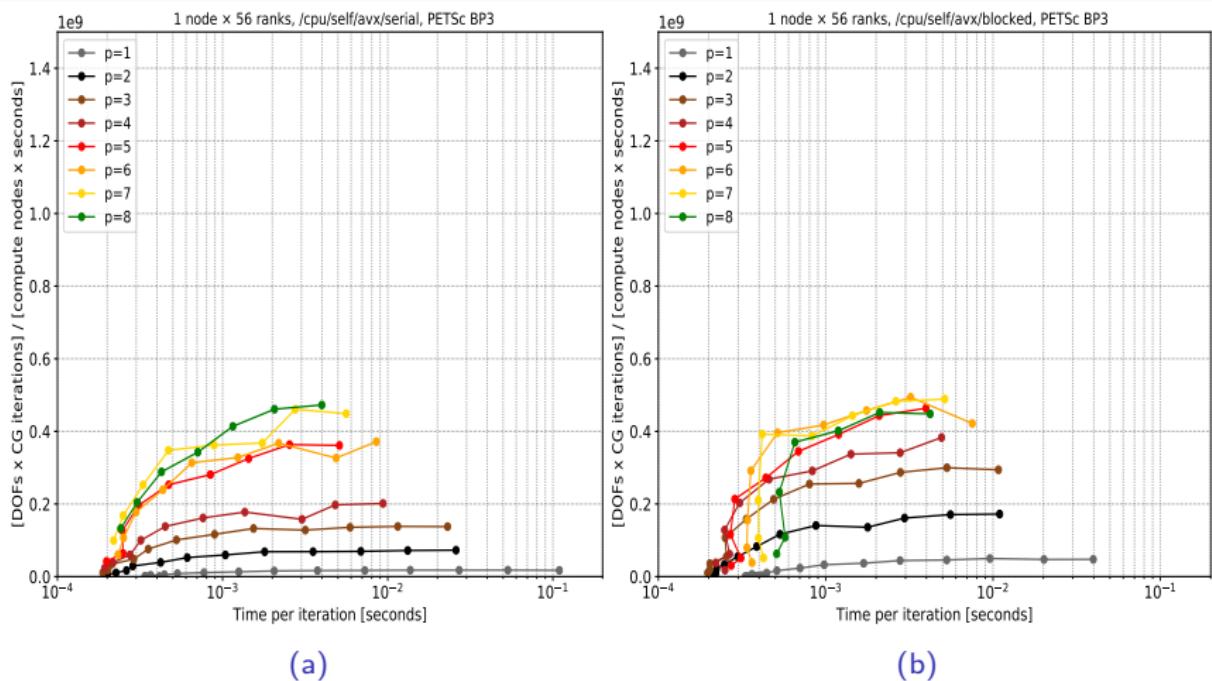
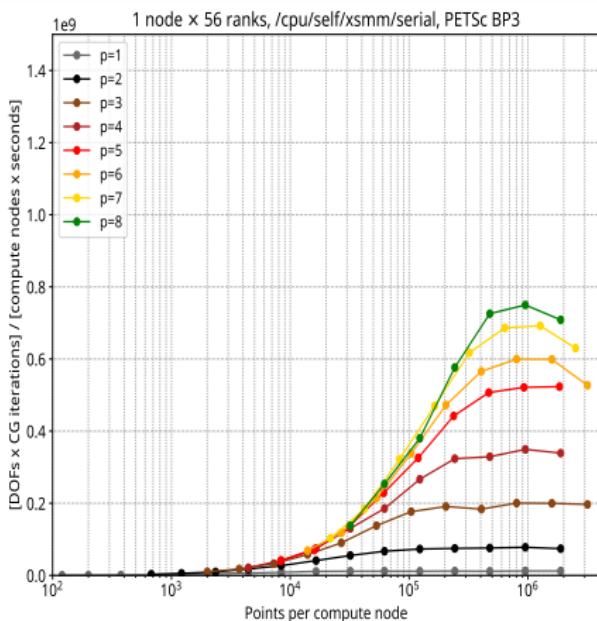
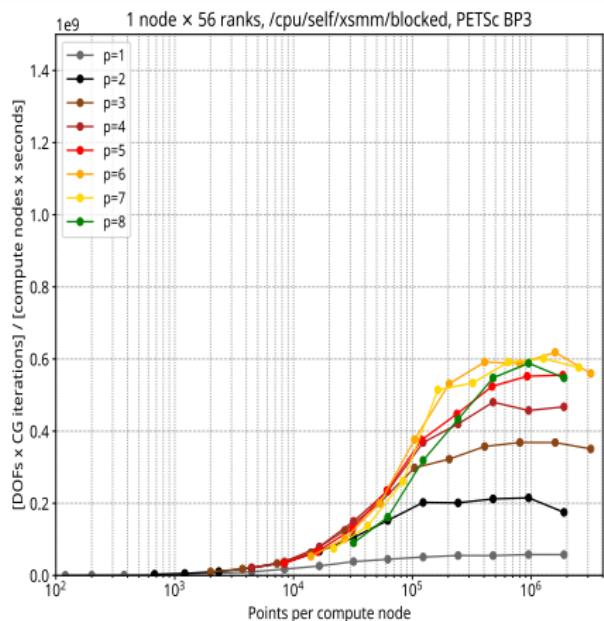


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Performance w. r. t. problem size on Skylake: libXSMM



(a)



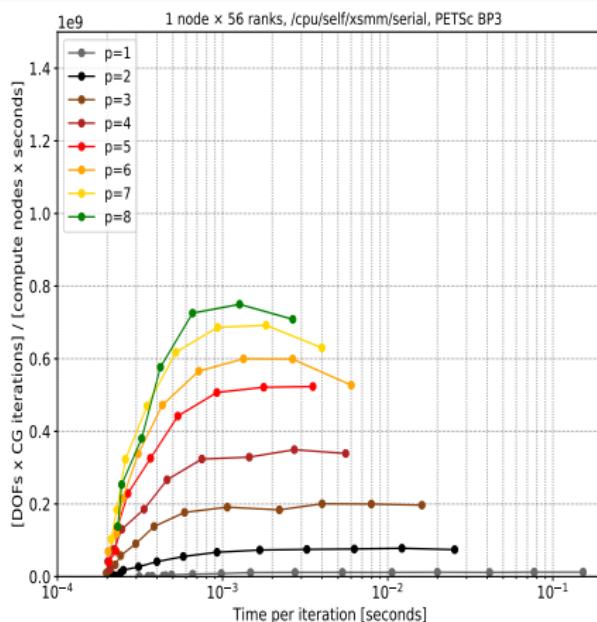
(b)

Figure: Skylake (2x Intel Xeon Platinum 8180M CPU 2.50GHz) with gcc-8 compiler. Backends: in (a) libXSMM serial; in (b) libXSMM blocked ($q = P + 2$, $P = p + 1$)

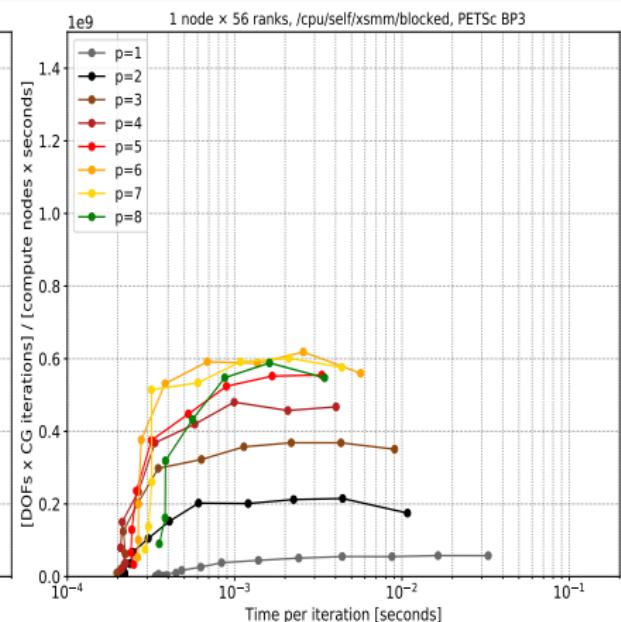


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Performance w. r. t. time on Skylake: libXSMM



(a)



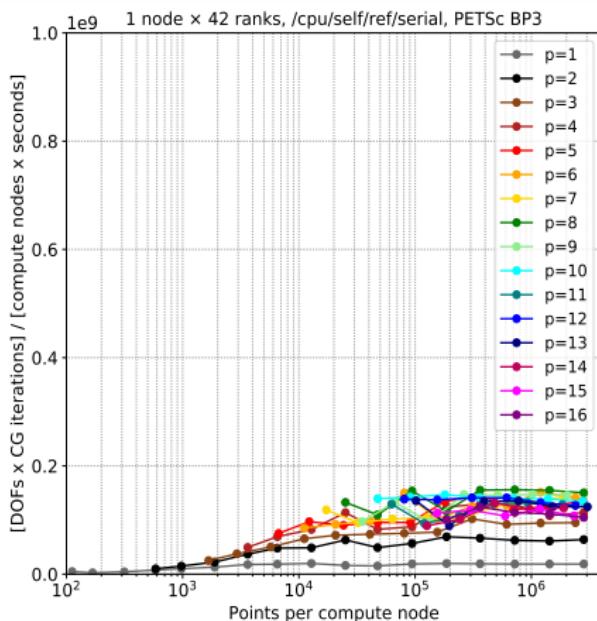
(b)

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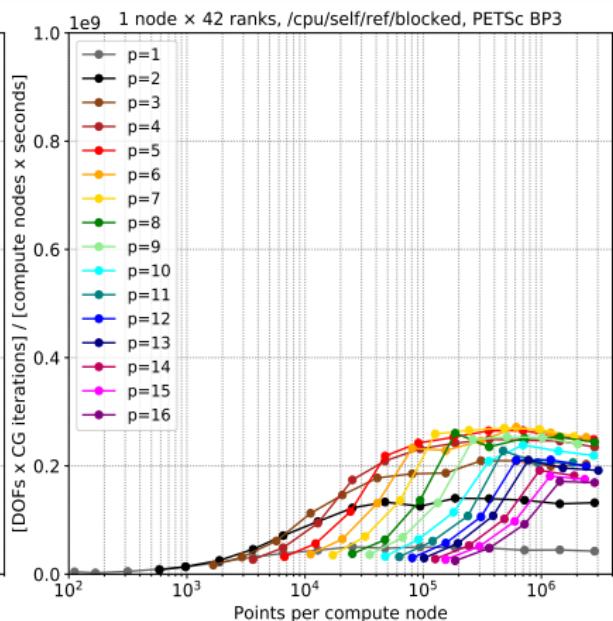


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Performance w. r. t. size on Summit: ref (non optimized)



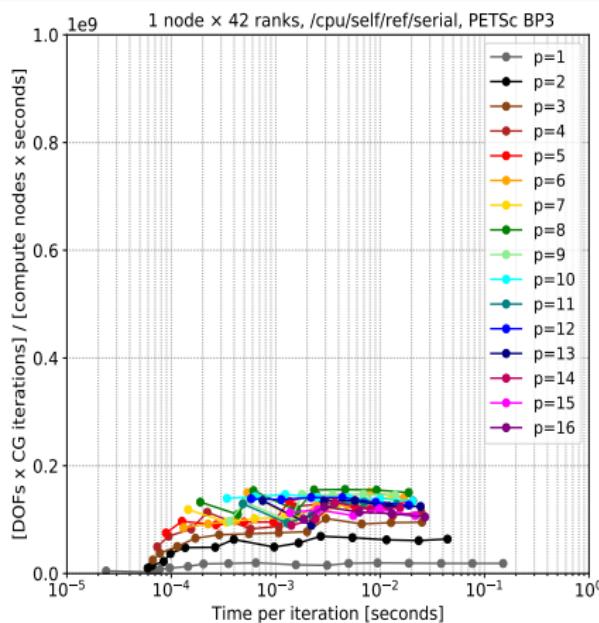
(a)



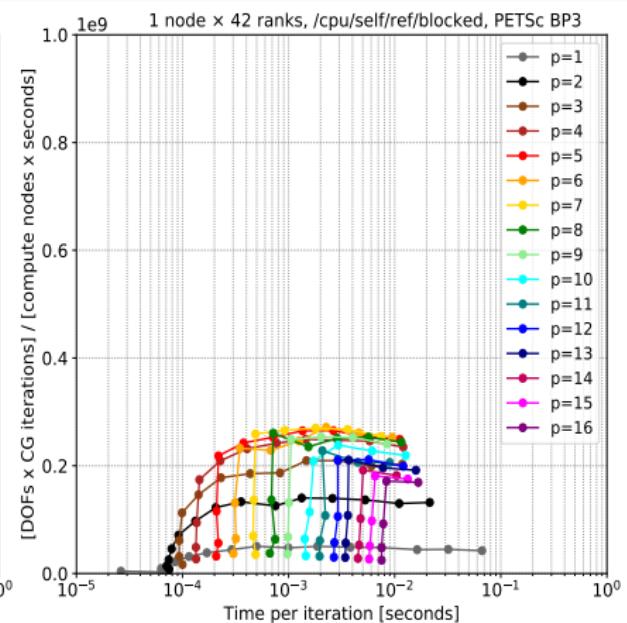
(b)

Figure: OLCF Summit (2x IBM POWER9) with gcc-9 compiler. Backends: in (a) ref serial; in (b) ref blocked ($q = P + 2$, $P = p + 1$)

Performance w. r. t. time on Summit: ref (non optimized)



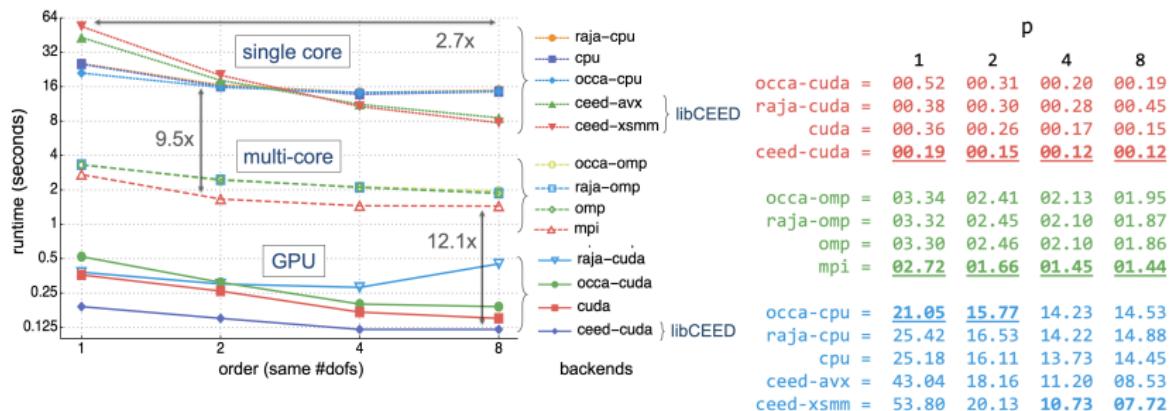
(a)



(b)

Figure: OLCF Summit (2x IBM POWER9) with gcc-9 compiler. Backends: in (a) ref serial; in (b) ref blocked ($q = P + 2$, $P = p + 1$)

Preliminary GPU results: MFEM + libCEED



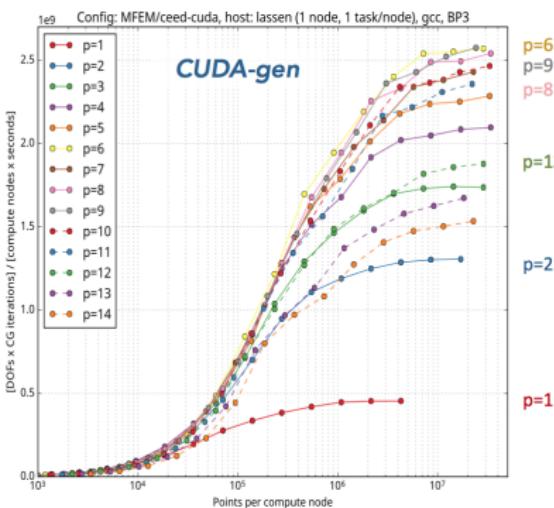
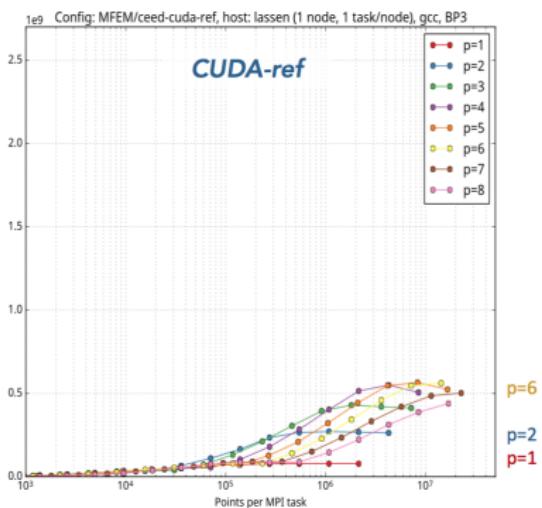
single-GPU, multi-core CPU, and single-core CPU,
for 1.3 millions DOFs in 2D.

Results by Yohann Dudouit on a Linux desktop with a Quadro GV100 GPU,
sm_70, CUDA 10.1, and Intel Xeon Gold 6130 CPU @ 2.10GHz.



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Preliminary GPU results: MFEM + libCEED (cont'ed)



Results by Yohann Dudouit on Lassen (LLNL): CUDA-ref (left) and CUDA-gen (right) backends performance for 3D BP3 on a NVIDIA V100 GPU.



Towards a libCEED miniapp: a Navier-Stokes solver

Compressible Navier-Stokes equations in conservation form:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \mathbf{u} = 0, \quad (1a)$$

$$\frac{\partial \mathbf{u}}{\partial t} + \nabla \cdot \left(\frac{\mathbf{u} \otimes \mathbf{u}}{\rho} + P \mathbf{I}_3 \right) + \rho g \mathbf{k} = \nabla \cdot \boldsymbol{\sigma}, \quad (1b)$$

$$\frac{\partial E}{\partial t} + \nabla \cdot \left(\frac{(E + P)\mathbf{u}}{\rho} \right) = \nabla \cdot (\mathbf{u} \cdot \boldsymbol{\sigma} + k \nabla T), \quad (1c)$$



Towards a libCEED miniapp: a Navier-Stokes solver

Compressible Navier-Stokes equations in conservation form:

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$$\frac{\partial E}{\partial t} + \nabla \cdot \left(\frac{(E + P)\mathbf{u}}{\rho} \right) = \nabla \cdot (\mathbf{u} \cdot \boldsymbol{\sigma} + k \nabla T), \quad (1c)$$

where $\boldsymbol{\sigma} = \mu(\nabla \mathbf{u} + (\nabla \mathbf{u})^T + \lambda(\nabla \cdot \mathbf{u})\mathbf{I}_3)$, and



Towards a libCEED miniapp: a Navier-Stokes solver

Compressible Navier-Stokes equations in conservation form:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \mathbf{u} = 0, \quad (1a)$$

$$\frac{\partial \mathbf{u}}{\partial t} + \nabla \cdot \left(\frac{\mathbf{u} \otimes \mathbf{u}}{\rho} + P \mathbf{I}_3 \right) + \rho g \mathbf{k} = \nabla \cdot \boldsymbol{\sigma}, \quad (1b)$$

$$\frac{\partial E}{\partial t} + \nabla \cdot \left(\frac{(E + P)\mathbf{u}}{\rho} \right) = \nabla \cdot (\mathbf{u} \cdot \boldsymbol{\sigma} + k \nabla T), \quad (1c)$$

where $\boldsymbol{\sigma} = \mu(\nabla \mathbf{u} + (\nabla \mathbf{u})^T + \lambda(\nabla \cdot \mathbf{u})\mathbf{I}_3)$, and

$$(c_p/c_v - 1)(E - \mathbf{u} \cdot \mathbf{u}/(2\rho) - \rho gz) = P \quad \leftarrow \text{pressure}$$

μ \leftarrow dynamic viscosity

g \leftarrow gravitational acceleration

k \leftarrow thermal conductivity

λ \leftarrow Stokes hypothesis constant

c_p \leftarrow specific heat, constant pressure

c_v \leftarrow specific heat, constant volume



Vector form

The system (1) can be rewritten in vector form

$$\frac{\partial \mathbf{q}}{\partial t} + \nabla \cdot \mathbf{F}(\mathbf{q}) = \mathbf{S}(\mathbf{q}), \quad (2)$$

for the state variables

$$\mathbf{q} = \begin{pmatrix} \rho \\ \mathbf{u} \equiv \rho \mathbf{u} \\ E \equiv \rho e \end{pmatrix} \leftarrow \begin{array}{l} \text{volume mass density} \\ \text{momentum density} \\ \text{energy density} \end{array} \quad (3)$$



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where

$$\mathbf{F}(\mathbf{q}) = \begin{pmatrix} \mathbf{u} \\ (\mathbf{u} \otimes \mathbf{u})/\rho + P\mathbf{I}_3 - \boldsymbol{\sigma} \\ (E + P)\mathbf{u}/\rho - (\mathbf{u} \cdot \boldsymbol{\sigma} + k\nabla T) \end{pmatrix},$$

$$\mathbf{S}(\mathbf{q}) = - \begin{pmatrix} 0 \\ \rho g \hat{\mathbf{k}} \\ 0 \end{pmatrix}$$



Space discretization

We use high-order finite/spectral elements: high-order Lagrange polynomials over non-uniformly spaced nodes, $\{x_i\}_{i=0}^p$, the Legendre-Gauss-Lobatto (LGL) points (roots of the p^{th} -order Legendre polynomial P_p). We let

$$\mathbb{R}^3 \supset \Omega = \bigcup_{e=1}^{N_e} \Omega_e, \text{ with } N_e \text{ disjoint hexaedral elements.}$$

The physical coordinates are $x = (x, y, z) \in \Omega_e$, while the reference coords are $X = (X, Y, Z) \in \mathbf{I} = [-1, 1]^3$.



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Define the discrete solution

$$\mathbf{q}_N(x, t)^{(e)} = \sum_{k=1}^P \psi_k(x) q_k^{(e)} \quad (4)$$

with P the number of nodes in the element e .

We use tensor-product bases $\psi_{kji} = h_i(X)h_j(Y)h_k(Z)$.



Strong and weak formulations

The strong form of (3):

$$\int_{\Omega} v \left(\frac{\partial q_N}{\partial t} + \nabla \cdot F(q_N) \right) d\Omega = \int_{\Omega} v S(q_N) d\Omega, \quad \forall v \in V_p \quad (5)$$

with $V_p = \{v \in H^1(\Omega_e) | v \in P_p(I), e = 1, \dots, N_e\}$.

Weak form:

$$\begin{aligned} & \int_{\Omega} v \frac{\partial q_N}{\partial t} d\Omega + \int_{\Gamma} v \hat{n} \cdot F(q_N) d\Omega - \int_{\Omega} \nabla v \cdot F(q_N) d\Omega = \\ & \int_{\Omega} v S(q_N) d\Omega, \quad \forall v \in V_p \end{aligned} \quad (6)$$



Strong and weak formulations

The strong form of (3):

$$\int_{\Omega} v \left(\frac{\partial \mathbf{q}_N}{\partial t} + \nabla \cdot \mathbf{F}(\mathbf{q}_N) \right) d\Omega = \int_{\Omega} v \mathbf{S}(\mathbf{q}_N) d\Omega, \quad \forall v \in \mathcal{V}_p \quad (5)$$

with $\mathcal{V}_p = \{v \in H^1(\Omega_e) | v \in P_p(I), e = 1, \dots, N_e\}$.

Weak form:

$$\begin{aligned} & \int_{\Omega} v \frac{\partial \mathbf{q}_N}{\partial t} d\Omega + \int_{\Gamma} v \hat{\mathbf{n}} \cdot \mathbf{F}(\mathbf{q}_N) d\Omega - \int_{\Omega} \nabla v \cdot \mathbf{F}(\mathbf{q}_N) d\Omega = \\ & \int_{\Omega} v \mathbf{S}(\mathbf{q}_N) d\Omega, \quad \forall v \in \mathcal{V}_p \end{aligned} \quad (6)$$

For the Time Discretization we use an explicit formulation

$$\frac{\mathbf{q}_N^{n+1} - \mathbf{q}_N^n}{\Delta t} = -[\nabla \cdot \mathbf{F}(\mathbf{q}_N)]^n + [\mathbf{S}(\mathbf{q}_N)]^n, \quad (7)$$

solved with the adaptive Runge-Kutta-Fehlberg (RKF4-5) method



A very simple example: The advection equation

We analyze the transport of total energy

$$\frac{\partial E}{\partial t} + \nabla \cdot (\mathbf{u}E) = 0, \quad (8)$$

with \mathbf{u} a uniform circular motion. BCs: no-slip and non-penetration for \mathbf{u} ,
no-flux for E .

order:
 $p = 6$
 $\Omega =$
 $[0, 2000]^3$ m
elem.
resolution:
250 m
Nodes: 117649



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Application example: Density current

A cold air bubble drops by convection in a neutrally stratified atmosphere.

Its initial condition is defined in terms of the Exner pressure, $\pi(x, t)$, and potential temperature, $\theta(x, t)$, that relate to the state variables via

$$\rho = \frac{P_0}{(c_p - c_v)\theta(x, t)} \pi(x, t)^{\frac{c_v}{c_p - c_v}}, \quad (9a)$$

$$e = c_v\theta(x, t)\pi(x, t) + \mathbf{u} \cdot \mathbf{u}/2 + gz, \quad (9b)$$

where P_0 is the atmospheric pressure.

BCs: no-slip and non-penetration for \mathbf{u} , no-flux for mass and energy densities.



Density current

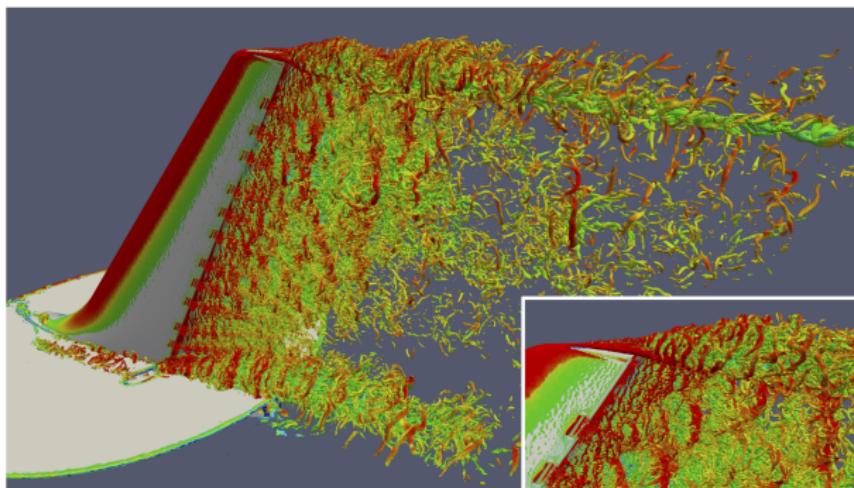
order: $p = 10$, $\Omega = [0, 6000]^2 \text{ m} \times [0, 3000] \text{ m}$, elem. resolution: 500 m,
Nodes: 893101



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Some work in progress

In collaboration with PHASTA (FastMath) we have worked on libCEED's integration.



[Ref: phasta.scigap.org]



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Stabilization methods

We are adding Streamline Upwind (SU) and Streamline Upwind/Petrov-Galerkin (SUPG) stabilization methods to our Navier-Stokes example.

For the advection case:

Not stabilized version.

Stabilized version.



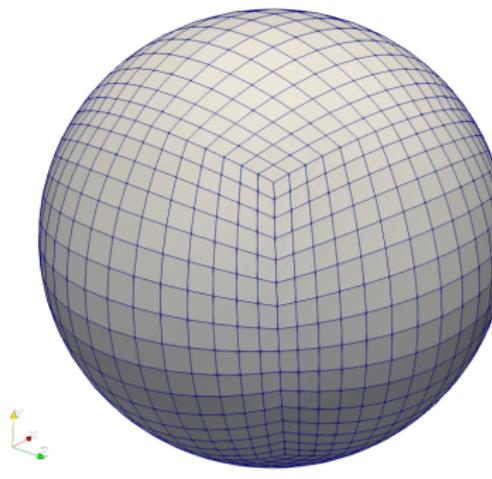
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Some other work in progress

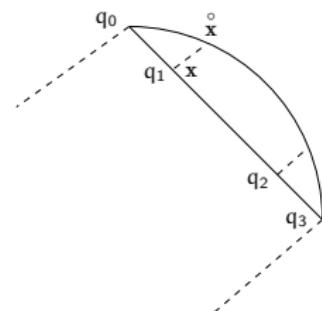
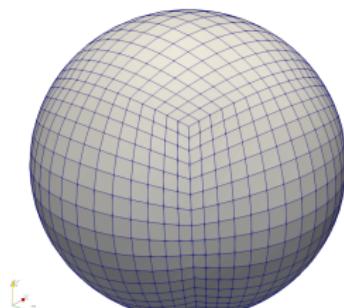
Converted BP1 (Mass operator) & BP3 (Poisson's equation) on the cubed-sphere as a prototype for shallow-water equations solver

$$\frac{\partial \mathbf{u}}{\partial t} = -(\omega + f)\hat{k} \times \mathbf{u} - \nabla \left(\frac{1}{2}|\mathbf{u}|^2 + g(h + h_s) \right) \quad (10a)$$

$$\frac{\partial h}{\partial t} = -\nabla \cdot (h_0 + h)\mathbf{u} \quad (10b)$$



Problems on the sphere



Transform $\overset{\circ}{\mathbf{x}} = (\overset{\circ}{x}, \overset{\circ}{y}, \overset{\circ}{z})$ on the sphere \hookrightarrow
 $\mathbf{x} = (x, y, z)$ on the discrete surface \hookrightarrow
 $\mathbf{X} = (X, Y) \in \mathbf{I} = [-1, 1]^2$

$$\frac{\partial \overset{\circ}{\mathbf{x}}}{\partial \mathbf{X} (3 \times 2)} = \frac{\partial \overset{\circ}{\mathbf{x}}}{\partial \mathbf{x} (3 \times 3)} \frac{\partial \mathbf{x}}{\partial \mathbf{X} (3 \times 2)}$$

$$|J| = \left| \text{col}_1 \left(\frac{\partial \overset{\circ}{\mathbf{x}}}{\partial \mathbf{X}} \right) \times \text{col}_2 \left(\frac{\partial \overset{\circ}{\mathbf{x}}}{\partial \mathbf{X}} \right) \right|$$



Conclusions

- We have showed libCEED's performance portability on several architectures, when integrated with PETSc and MFEM
- We have demonstrated the use of libCEED with PETSc for the numerical high-order solutions of
 - Advection equation
 - Full compressible Navier-Stokes equations
- We welcome contributors and users
<https://github.com/CEED/libCEED>



Outlook

Ongoing and future work:

- Implicit-explicit (IMEX) and fully implicit time discretizations for:
 - NS solver & SWE solver
- SU/SUPG stabilized version of the NS solver
- BDDC and FDM preconditioners
- Algorithmic differentiation of Q-functions
- Ongoing work on CUDA and HIP optimizations
- Our (first!) user manual and educational Jupyter-notebook tutorials



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Thank you!



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Tensor contractions

Let $\{x_i\}_{i=0}^p$ denote the LGL nodes with the corresponding interpolants $\{\psi_i^p\}_{i=0}^p$. Choose a quadrature rule with nodes $\{q_i^Q\}_{i=0}^Q$ and weights $\{w_i^Q\}$. The basis evaluation, derivative, and integration matrices are $B_{ij}^{Qp} = \psi_j^p(q_i^Q)$, $D_{ij}^{Qp} = \partial_x \psi_j^p(q_i^Q)$, and $W_{ij}^Q = w_i^Q \delta_{ij}$. In 3D:

$$\mathbf{B} = \mathbf{B} \otimes \mathbf{B} \otimes \mathbf{B} \quad (11)$$

$$\mathbf{D}_0 = \mathbf{D} \otimes \mathbf{B} \otimes \mathbf{B} \quad (12)$$

$$\mathbf{D}_1 = \mathbf{B} \otimes \mathbf{D} \otimes \mathbf{B} \quad (13)$$

$$\mathbf{D}_2 = \mathbf{B} \otimes \mathbf{B} \otimes \mathbf{D} \quad (14)$$

$$\mathbf{W} = \mathbf{W} \otimes \mathbf{W} \otimes \mathbf{W} \quad (15)$$

These tensor-product operations cost $2(p^3 Q + p^2 Q^2 + p Q^3)$ and touch only $O(p^3 + Q^3)$ memory. In the spectral element method, when the same LGL points are reused for quadrature (i.e., a collocated method with $Q = p + 1$), then $\mathbf{B} = \mathbf{I}$ and \mathbf{D} reduces to $O(p^4)$.